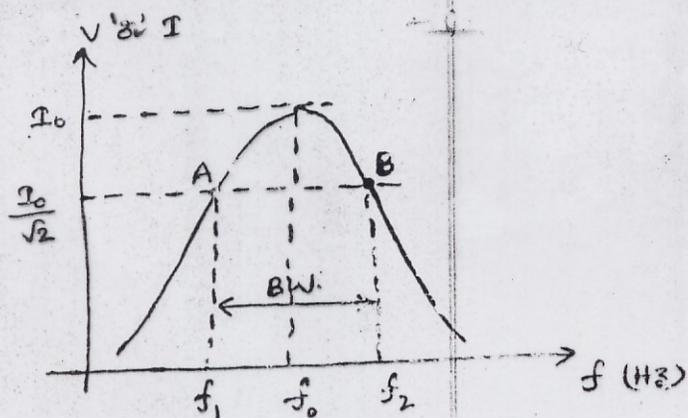


Bandwidth of a series Resonant ckt :-

The bandwidth of any system is the range of frequencies for which the current or output voltage is equal to 70.7% of its value at the resonant frequency.



Bandwidth (BW) is also defined as the frequency betw difference between f_2 and f_1 .

$$\therefore BW = f_2 - f_1$$

(or)

$$\Delta f = f_2 - f_1$$

(or)

$$\Delta \omega = \omega_2 - \omega_1$$

f_1 : lower cut-off frequency.

f_2 : upper cut-off frequency.

power dissipated in the ckt at the frequency f_1 is,

$$P_A = I_A^2 R = \left(\frac{I_0}{\sqrt{2}}\right)^2 R = \frac{I_0^2 R}{2}$$

power dissipated in the ckt. at

frequency f_2 is,

$$P_B = I_B^2 R = \left(\frac{I_0}{\sqrt{2}}\right)^2 R = \frac{I_0^2 R}{2}$$

$$\therefore P_A = P_B = \frac{I_0^2 R}{2}$$

where $I_0^2 R = P_0$.

$$\therefore P_A = P_B = \frac{P_0}{2}$$

For this reason, the two points A & B on the frequency response curve are known as "half power points".

The Decibel power response :-

$$\text{power ratio} = 10 \log_{10} \left[\frac{\text{power at half power points}}{\text{power at resonance}} \right]$$

$$= 10 \log_{10} \left[\frac{P_0/2}{P_0} \right]$$

$$= 10 \log_{10} \left[\frac{1}{2} \right]$$

$$= -3 \text{ dB.}$$

Hence the half power points are also called '3dB' points. At these points, the power is '3dB' less than at resonance.

Quality Factor (Q-factor) &v

Figure of merit :-

Q-factor is defined as,

$$Q = \frac{\text{max. energy stored/cycle}}{\text{Energy dissipated/cycle}}$$

For an inductor :-

$$\text{max energy stored} = \frac{1}{2} L I_m^2$$

$$\text{energy dissipated/cycle} = \left(\frac{I_m}{\sqrt{2}}\right)^2 RT$$

$$= \frac{I_m^2 RT}{2}$$

$$= \frac{I_m^2 R}{2} \quad (1)$$

$$\therefore Q = 2\pi \left[\frac{\frac{1}{2} L I_m^2}{\frac{I_m^2 R}{2} \left(\frac{1}{f}\right)} \right]$$

$$Q = \frac{2\pi f L}{R}$$

$$Q = \frac{\omega L}{R}$$

For a capacitor :-

$$\begin{aligned} \text{Max. energy stored} &= \frac{1}{2} C V_m^2 \\ &= \frac{1}{2} C \left[I_m X_C \right]^2 \\ &= \frac{1}{2} C \left[\frac{I_m}{\omega^2 C^2} \right] \\ &= \frac{1}{2} \frac{I_m^2}{\omega^2 C} \end{aligned}$$

$$\begin{aligned} \therefore \text{Energy dissipated/cycle} &= \left(\frac{I_m}{\sqrt{2}} \right)^2 R \times T \\ &= \frac{I_m^2 R T}{2} \\ &= \frac{I_m^2 R}{2} \left(\frac{1}{f} \right) \end{aligned}$$

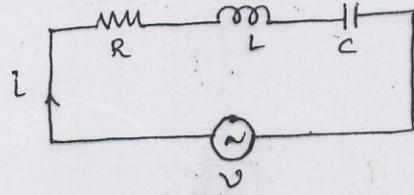
$$\therefore Q = 2\pi \left[\frac{\frac{1}{2} \frac{I_m^2}{\omega^2 C}}{\frac{I_m^2 R}{2} \left(\frac{1}{f} \right)} \right]$$

$$\therefore Q = \frac{2\pi f}{\omega^2 R C} = \frac{2\pi f}{(2\pi f)^2 R C} = \frac{1}{2\pi f R C}$$

$$Q = \frac{1}{\omega R C}$$

$$\therefore \text{Q factor 'Q'} = \frac{\omega L}{R} = \frac{1}{\omega R C}$$

Band Width Interms of ckt parameters :-



$$\begin{aligned} \text{Impedance } Z &= R + j(X_L - X_C) \\ &= R + j\left(\omega L - \frac{1}{\omega C}\right) \end{aligned}$$

$$|Z| = \sqrt{R^2 + \left(\omega L - \frac{1}{\omega C}\right)^2}$$

$$\therefore \text{current } I = \frac{V}{|Z|}$$

$$I = \frac{V}{\sqrt{R^2 + \left(\omega L - \frac{1}{\omega C}\right)^2}} \rightarrow \textcircled{1}$$

At half power points,

$$I_0 = \frac{V}{Z} \quad \text{But } Z=R$$

$$\therefore I_0 = \frac{V}{R}$$

$$\text{and current } I = \frac{I_0}{\sqrt{2}}$$

$$\therefore I = \frac{V/R}{\sqrt{2}} = \frac{V}{\sqrt{2}R} \rightarrow \textcircled{2}$$

\therefore Equate eq ① & eq ②;

$$\frac{V}{\sqrt{2}R} = \frac{V}{\sqrt{R^2 + \left(\omega L - \frac{1}{\omega C}\right)^2}}$$

$$\sqrt{2}R = \sqrt{R^2 + \left(\omega L - \frac{1}{\omega C}\right)^2}$$

Squaring on both sides,

$$2R^2 = R^2 + \left(\omega L - \frac{1}{\omega C}\right)^2$$

$$2R^2 - R^2 = \left(\omega L - \frac{1}{\omega C}\right)^2$$

$$\therefore \omega L - \frac{1}{\omega C} = \pm R \rightarrow \textcircled{3}$$

We know that $Z = R + j(\omega L - \frac{1}{\omega C})$

$$\therefore Z = R + jR$$

$$\therefore \boxed{Z = R(1 + j)} \rightarrow \textcircled{4}$$

The above equation gives the impedance at half power frequencies.

Let ω_1 & ω_2 be the lower half power and upper half power frequencies respectively.

At lower half-power frequency,

eq (3) becomes,

$$\omega_1 L - \frac{1}{\omega_1 C} = -R \quad (\because '-' \text{ for capacitive nature})$$

$$\frac{\omega_1^2 LC - 1}{\omega_1 C} = -R$$

$$\omega_1^2 LC + \omega_1 RC - 1 = 0.$$

$$\therefore \omega_1 = \frac{-(RC) \pm \sqrt{(RC)^2 - 4(LC)(-1)}}{2LC}$$

Consider the +ve value for ' ω_1 ';

$$\boxed{\omega_1 = -\frac{R}{2L} + \sqrt{\left(\frac{R}{2L}\right)^2 + \frac{1}{LC}} \rightarrow \textcircled{5}}$$

Similarly, at higher half-power freq,

eq (3) becomes,

$$\omega_2 L - \frac{1}{\omega_2 C} = +R \quad (\because '+' \text{ for inductive nature})$$

$$\omega_2^2 LC - \omega_2 RC - 1 = 0$$

$$\therefore \omega_2 = \frac{+(RC) \pm \sqrt{(RC)^2 - 4(LC)(-1)}}{(2)LC}$$

Consider the +ve value for ' ω_2 ';

$$\boxed{\omega_2 = \frac{R}{2L} + \sqrt{\left(\frac{R}{2L}\right)^2 + \frac{1}{LC}} \rightarrow \textcircled{6}}$$

We know that,

$$BW, \Delta f = f_2 - f_1 \quad (\text{or})$$

$$\Delta \omega = \omega_2 - \omega_1$$

$$\therefore BW = \text{eq (6)} - \text{eq (5)}$$

$$\Delta \omega = 2 \left[\frac{R}{2L} \right]$$

$$\therefore \boxed{\Delta \omega = \frac{R}{L}}$$

(OR)

$$\boxed{\Delta f = \frac{R}{2\pi L}}$$

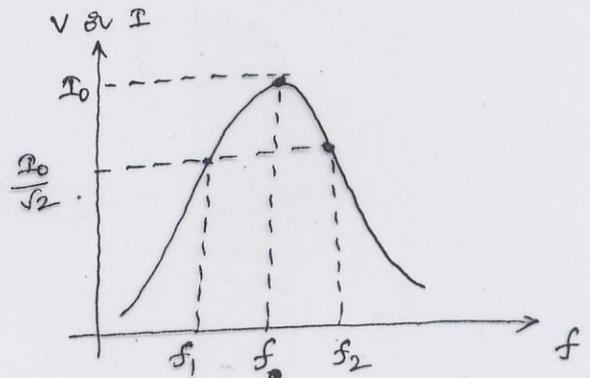
\therefore The B.W. of a series RLC ckt depends upon the ' $\frac{R}{L}$ ' ratio and independent of 'C'.

Special case:

We know that,

$$\omega_1 = -\frac{R}{2L} + \sqrt{\left(\frac{R}{2L}\right)^2 + \frac{1}{LC}} \rightarrow \textcircled{1}$$

$$\omega_2 = \frac{R}{2L} + \sqrt{\left(\frac{R}{2L}\right)^2 + \frac{1}{LC}} \rightarrow \textcircled{2}$$



$$\omega_1 \omega_2 = \left[\sqrt{\left(\frac{R}{2L}\right)^2 + \frac{1}{LC}} \right]^2 - \left[\frac{R}{2L} \right]^2$$

$$\therefore \omega_1 \omega_2 = \left(\frac{1}{LC} \right) \rightarrow \textcircled{3}$$

We know that $f_0 = \frac{1}{2\pi\sqrt{LC}}$

$$\therefore \omega_b^2 = \frac{1}{LC}$$

eq (3) can be written as,

$$\omega_1 \omega_2 = \omega_0^2$$

OR

$$\omega_b = \sqrt{\omega_1 \omega_2}$$

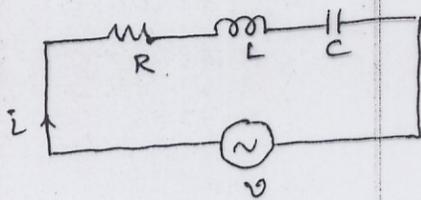
Resonant frequency in terms of half power frequency is given by,

$$2\pi f_0 = \sqrt{2\pi f_1 \cdot 2\pi f_2}$$

$$* f_0 = \sqrt{f_1 f_2} *$$

Hence, the "resonant frequency" is the geometric mean of the two half-power frequencies.

Relationship between Bandwidth, Resonant frequency & Quality factor :-



For a series R-L-C circuit,

$$\begin{aligned} Z &= R + j(X_L - X_C) \\ &= R + j(\omega L - \frac{1}{\omega C}) \\ &= R \left[1 + j \left(\frac{\omega L}{R} - \frac{1}{\omega RC} \right) \right] \end{aligned}$$

$$= R \left[1 + j \left\{ \frac{\omega L}{R} \times \frac{\omega_0}{\omega_0} - \frac{1}{\omega RC} \times \frac{\omega_0}{\omega_0} \right\} \right]$$

$$\therefore Z = R \left[1 + j \left(\frac{\omega_0 L}{R} \cdot \frac{\omega}{\omega_0} - \frac{1}{\omega_0 RC} \cdot \frac{\omega_0}{\omega} \right) \right]$$

We know that, at resonance,

$$\begin{aligned} \therefore Z &= R \left[1 + j \left(Q_0 \cdot \frac{\omega}{\omega_0} - Q_0 \cdot \frac{\omega_0}{\omega} \right) \right] \\ &= R \left[1 + j Q_0 \left(\frac{\omega}{\omega_0} - \frac{\omega_0}{\omega} \right) \right] \rightarrow (1) \end{aligned}$$

$$\text{Let } \delta = \frac{\omega - \omega_0}{\omega_0} = \frac{\omega}{\omega_0} - 1$$

" δ " is called as "fractional freq variation" or "fractional detuning factor" or "per unit freq deviation".

$$\therefore \delta = \frac{\omega}{\omega_0} - 1 \Rightarrow (1 + \delta) = \frac{\omega}{\omega_0}$$

From eq (1);

$$\begin{aligned} Z &= R \left[1 + j Q_0 \left[(1 + \delta) - \frac{1}{(1 + \delta)} \right] \right] \\ &= R \left[1 + j Q_0 \left[\frac{\delta(2 + \delta)}{(1 + \delta)} \right] \right] \end{aligned}$$

Let $\delta \ll 1$; then

$$Z = R \left[1 + j 2 Q_0 \delta \right] \rightarrow (2)$$

The above equation represents the relationship between 'R' and 'delta'.

At half power frequencies, the impedance is given by,

$$Z = R(1 \pm j1) \rightarrow (3)$$

equating (2) & (3);

$$2 Q_0 \delta = \pm 1$$

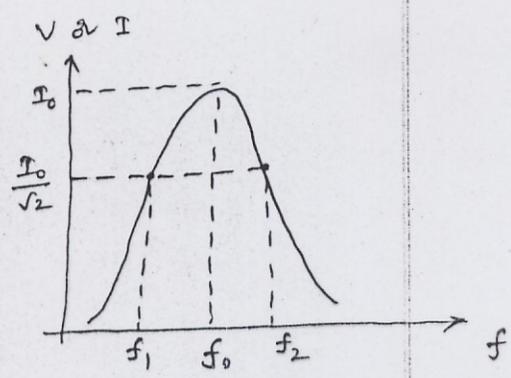
$$\delta = \pm \frac{1}{2 Q_0} \rightarrow (4)$$

At ω_1 ;

$$\delta = \frac{\omega - \omega_0}{\omega_0} = \frac{\omega_1 - \omega_0}{\omega_0}$$

At ω_2 ;

$$\delta = \frac{\omega - \omega_0}{\omega_0} = \frac{\omega_2 - \omega_0}{\omega_0}$$



$$\therefore \frac{\omega_1 - \omega_0}{\omega_0} = \frac{f_1 - f_0}{f_0} = -\frac{1}{2Q_0} \rightarrow \text{(A)}$$

$$\therefore \frac{\omega_2 - \omega_0}{\omega_0} = \frac{f_2 - f_0}{f_0} = +\frac{1}{2Q_0} \rightarrow \text{(B)}$$

(B) - (A);

$$\frac{1}{2Q_0} + \frac{1}{2Q_0} = \frac{f_2 - f_0}{f_0} - \left(-\frac{f_1 - f_0}{f_0} \right)$$

$$\frac{1}{Q_0} = \frac{f_2 - f_1}{f_0}$$

$$Q_0 = \frac{f_0}{f_2 - f_1}$$

$$Q_0 = \frac{f_0}{\text{BW}}$$

\therefore Quality factor = $\frac{\text{Resonant frequency}}{\text{Bandwidth}}$.

$$Q_0 \propto \frac{1}{\text{BW}} ; Q_0 \propto f_0$$

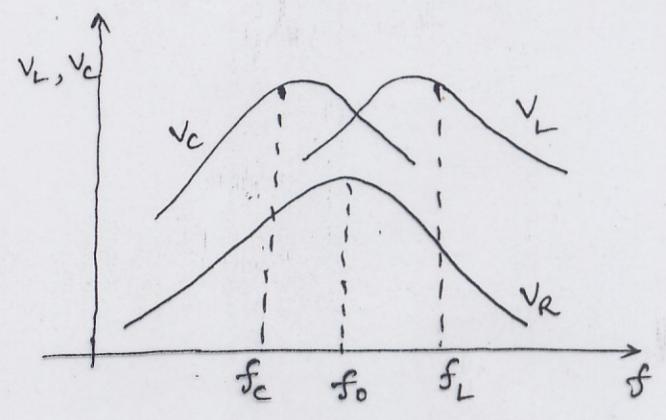
Note:

"Selectivity" of a resonant ckt is its ability to discriminate between signals of desired and undesired frequencies.

$$\text{Selectivity} = S = \frac{\omega_0 L}{R} = Q$$

Thus for a series RLC ckt, the selectivity is equal to the quality factor.

Variation of current & voltage with frequency :-
 (Frequency at which V_C & V_L are max)



for a series R-L-C ckt,

$$Z = R + j(X_L - X_C)$$

$$Z = R + j\left(\omega L - \frac{1}{\omega C}\right)$$

$$I = \frac{V}{|Z|} = \frac{V}{\sqrt{R^2 + \left(\omega L - \frac{1}{\omega C}\right)^2}}$$

(At resonance; $Z=R$; then $I_0 = \frac{V}{R}$. Hence current is max, at resonance).

voltage across capacitor is

$$V_C = I X_C = I \left(\frac{1}{\omega C}\right)$$

$$\therefore V_C = \left[\frac{V}{\sqrt{R^2 + \left(\omega L - \frac{1}{\omega C}\right)^2}} \right] \left(\frac{1}{\omega C}\right)$$

$$V_C^2 = \frac{V^2}{R^2 + \left(\omega L - \frac{1}{\omega C}\right)^2} \cdot \frac{1}{\omega^2 C^2}$$

$$V_C^2 = \frac{V^2}{\omega^2 C^2 \left(R^2 + \left(\frac{\omega^2 LC - 1}{\omega C}\right)^2 \right)}$$

$$= \frac{V^2}{\omega^2 C^2 \left[R^2 + \frac{(\omega^2 LC - 1)^2}{\omega^2 C^2} \right]}$$

$$= \frac{V^2}{R^2 + (\omega^2 LC - 1)^2}$$

$$\therefore V_c^r = \frac{V^r}{R^r \omega^r C^r + (\omega^r L C^r - 1)^r} \rightarrow \textcircled{1}$$

The frequency at which V_c^r is max, can be obtained by equating $\frac{dV_c^r}{d\omega} = 0$.

$$\frac{d}{d\omega} [V_c^r] = 0 \rightarrow \frac{V^r [2(\omega) [R^r C^r] + 2 [\omega^r L C^r - 1] \cdot (2\omega) (LC)]}{(\omega^r R^r C^r + [\omega^r L C^r - 1]^r)^2} = 0$$

$$\therefore V^r [2\omega R^r C^r + 2(\omega^r L C^r - 1) 2\omega LC] = 0$$

$V \neq 0$; But

$$2\omega R^r C^r + 2(\omega^r L C^r - 1) 2\omega LC = 0$$

$$2\omega R^r C^r + 4\omega^3 L^2 C^r - 4\omega LC = 0$$

$$2\omega C [R^r C^r + 2\omega^2 L^2 C^r - 2L] = 0$$

$$2\omega^2 L^2 C^r + R^r C^r - 2L = 0$$

$$2\omega^2 L^2 C^r = 2L - R^r C^r$$

$$\omega^r = \frac{2L}{2L^2 C^r} - \frac{R^r C^r}{2L^2 C^r}$$

put, $\omega = \omega_c$;

$$\omega_c^r = \frac{1}{LC} - \frac{R^r}{2L^2}$$

$$\omega_c = \sqrt{\frac{1}{LC} - \frac{R^r}{2L^2}} \text{ rad/sec.}$$

(or)

$$f_c = \frac{1}{2\pi} \sqrt{\frac{1}{LC} - \frac{R^r}{2L^2}} \rightarrow \textcircled{2}$$

Similarly;

voltage across the inductor, $V_L = I \cdot X_L$

$$\text{But } I = \frac{V}{\sqrt{R^r + (\omega L - \frac{1}{\omega C})^r}}$$

$$\therefore V_L = \left[\frac{V}{\sqrt{R^r + (\omega L - \frac{1}{\omega C})^r}} \right] \cdot (\omega L)$$

$$V_L^r = \frac{V^r \omega^r L^r}{R^r + (\omega L - \frac{1}{\omega C})^r}$$

$$= \frac{V^r \omega^r L^r}{R^r + \left(\frac{\omega^r L C^r - 1}{\omega C}\right)^r}$$

$$\therefore V_L^r = \frac{V^r \omega^r L^r}{\left[R^r \omega^r C^r + (\omega^r L C^r - 1)^r \right] \omega^r C^r}$$

$$\therefore V_L^r = \frac{V^r \omega^4 L^2 C^r}{R^r \omega^r C^r + (\omega^r L C^r - 1)^r} \rightarrow \textcircled{3}$$

The frequency at which V_L^r max, can be obtained by equating $\frac{dV_L^r}{d\omega} = 0$.

Then we will get

$$2\omega^3 L C^r - \omega^r R^r C^r - 2 = 0$$

$$\omega^r [2LC - C^r R^r] = 2$$

$$\omega^r = \frac{2}{2LC - C^r R^r} = \frac{1}{LC - \frac{C^r R^r}{2}}$$

put $\omega = \omega_L$

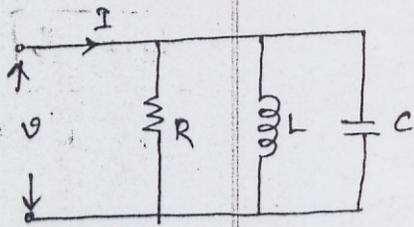
$$\omega_L = \frac{1}{\sqrt{LC - \frac{R^r C^r}{2}}} \text{ rad/sec.}$$

(or)

$$f_L = \frac{1}{2\pi} \left(\frac{1}{\sqrt{LC - \frac{R^r C^r}{2}}} \right) \rightarrow \textcircled{4}$$

Eq ② & ④ represents the expressions of the frequencies at which V_c^r & V_L^r are maximum.

Parallel Resonance for ideal ckt :-



At resonance, in parallel ckt, the net susceptance must be zero. For analyzing parallel ckt, admittance-concept is more convenient.

The admittances of each branch are:

$$Y_1 = \frac{1}{R} = G$$

$$Y_2 = \frac{1}{jX_L} = -j\left(\frac{1}{X_L}\right) = -jB_L$$

$$Y_3 = \frac{1}{-jX_C} = +j\left(\frac{1}{X_C}\right) = +jB_C$$

Total admittance $Y = Y_1 + Y_2 + Y_3$.

$$\therefore Y = G + j(B_C - B_L)$$

Condition for resonance is net susceptance should be zero.

$$\therefore B_C - B_L = 0$$

$$B_C = B_L$$

$$\frac{1}{X_C} = \frac{1}{X_L}$$

$$\omega_0 C = \frac{1}{\omega_0 L}$$

$$\omega_0^2 = \frac{1}{LC}$$

$$\omega_0 = \frac{1}{\sqrt{LC}} \text{ rad/sec.}$$

(OR)

$$f_0 = \frac{1}{2\pi\sqrt{LC}} \text{ Hz.}$$

Bandwidth & Q-factor for a parallel R-L-C circuit :-

For a parallel R-L-C ckt,

$$Y = G + j(B_C - B_L)$$

We know that,

$$\omega L - \frac{1}{\omega C} = \pm R \text{ (for series ckt)}$$

Similarly for parallel ckt,

$$\omega C - \frac{1}{\omega L} = \pm \frac{1}{R}$$

At $\omega = \omega_1$;

$$\omega_1 C - \frac{1}{\omega_1 L} = -\frac{1}{R}$$

$$\frac{\omega_1^2 LC - 1}{\omega_1 L} = -\frac{1}{R}$$

$$\omega_1^2 RLC - R = -\omega_1 L$$

$$\omega_1^2 RLC + \omega_1 L - R = 0.$$

$$\omega_1^2 + \frac{\omega_1}{RC} - \frac{1}{LC} = 0$$

$$\omega_1 = \frac{-\frac{1}{RC} \pm \sqrt{\left(\frac{1}{RC}\right)^2 - 4(1)\left(-\frac{1}{LC}\right)}}{2(1)} \quad (2) (1)$$

$$\therefore \omega_1 = \frac{-\frac{1}{2RC} + \sqrt{\left(\frac{1}{2RC}\right)^2 + \frac{1}{LC}}}{1} \rightarrow (1)$$

Similarly,

$$\omega_2 = \frac{\frac{1}{2RC} + \sqrt{\left(\frac{1}{2RC}\right)^2 + \frac{1}{LC}}}{1} \rightarrow (2)$$

$$B-W = \Delta\omega = \omega_2 - \omega_1$$

$$\therefore \Delta\omega = BW = \frac{1}{RC}$$

(OR)

$$\Delta f = \frac{1}{2\pi RC}$$

$$\therefore Q = \frac{\omega_0}{\Delta\omega} = \omega_0 RC$$

where $\omega_0 = \frac{1}{\sqrt{LC}}$

$$\therefore Q = \frac{\omega_0^2 RC}{\omega_0} = \left[\left(\frac{1}{LC} \right) RC \right] / \omega_0$$

$$\therefore Q = \frac{R}{\omega_0 L}$$

Resonant frequency for a practical parallel R-L-C circuit :-

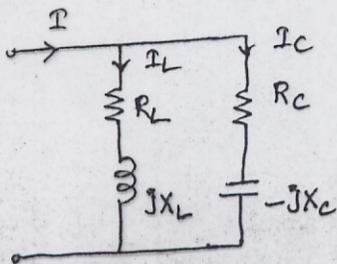


fig: (a)

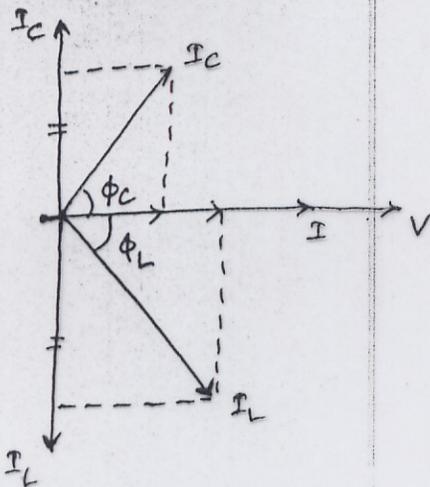


fig: (b)

The condition for resonance in the parallel ckt is $I_C \sin \phi_C = I_L \sin \phi_L$

From fig: (a)

$$Y = Y_L + Y_C$$

$$= \frac{1}{Z_L} + \frac{1}{Z_C}$$

$$\begin{aligned} & \frac{R_L + jX_L}{R_L^2 + X_L^2} + \frac{R_C - jX_C}{R_C^2 + X_C^2} \\ &= \left[\frac{R_L}{R_L^2 + X_L^2} + \frac{R_C}{R_C^2 + X_C^2} \right] + \\ & \quad j \left[\frac{X_C}{R_C^2 + X_C^2} - \frac{X_L}{R_L^2 + X_L^2} \right] \end{aligned}$$

At resonance, the reactive part of $Y=0$

$$\text{i.e. } \frac{X_C}{R_C^2 + X_C^2} - \frac{X_L}{R_L^2 + X_L^2} = 0$$

$$\therefore X_C [R_L^2 + X_L^2] - X_L [R_C^2 + X_C^2] = 0$$

$$\frac{1}{\omega_0 C} [R_L^2 + \omega_0^2 L^2] = \omega_0 L [R_C^2 + \frac{1}{\omega_0^2 C^2}]$$

$$R_L^2 + \omega_0^2 L^2 = \omega_0^2 LC [R_C^2 + \frac{1}{\omega_0^2 C^2}]$$

$$R_L^2 + \omega_0^2 L^2 = \omega_0^2 LC R_C^2 + \frac{L}{C}$$

$$\omega_0^2 L^2 - \omega_0^2 LC R_C^2 = \frac{L}{C} - R_L^2$$

$$\omega_0^2 [L^2 - R_C^2 LC] = \frac{L}{C} - R_L^2$$

$$\omega_0^2 LC \left[\frac{L}{C} - R_C^2 \right] = \frac{L}{C} - R_L^2$$

$$\omega_0^2 LC \left[R_C^2 - \frac{L}{C} \right] = R_L^2 - \frac{L}{C}$$

$$\omega_0^2 = \frac{1}{LC} \left[\frac{R_L^2 - \frac{L}{C}}{R_C^2 - \frac{L}{C}} \right]$$

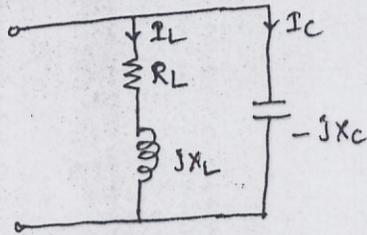
$$\omega_0 = \sqrt{\frac{1}{LC} \left[\frac{R_L^2 - \frac{L}{C}}{R_C^2 - \frac{L}{C}} \right]} \text{ rad/sec}$$

(00)

$$f_0 = \frac{1}{2\pi} \sqrt{\frac{1}{LC} \left(\frac{R_L^2 - \frac{L}{C}}{R_C^2 - \frac{L}{C}} \right)} \text{ Hz}$$

Special case :-

In a more practical ckt, the resistance in the capacitance branch is negligible. " $R_c = 0$," then



$$f_0 = \frac{1}{2\pi} \sqrt{\frac{1}{LC} \left(\frac{R_L^2 - L/C}{R_c^2 - L/C} \right)}$$

where $R_c = 0$;

$$f_0 = \frac{1}{2\pi} \sqrt{\frac{1}{LC} \left(\frac{R_L^2 - L/C}{-L/C} \right)}$$

$$\therefore f_0 = \frac{1}{2\pi} \sqrt{\frac{1}{LC} \left(1 - \frac{R_L^2 C}{L} \right)}$$

(OR)

$$f_0 = \frac{1}{2\pi} \sqrt{\frac{1}{LC} - \left(\frac{R_L}{L} \right)^2}$$

Impedance at Resonance :-

$$Y_0 = \frac{1}{Z_0}$$

$$Y_0 = \frac{R_L}{R_L^2 + \omega_0^2 L^2} \rightarrow \textcircled{1}$$

we know that,

$$\omega_0 = \sqrt{\frac{1}{LC} \left(1 - \frac{R_L^2 C}{L} \right)}$$

$$\omega_0^2 = \frac{1}{LC} \left(1 - \frac{R_L^2 C}{L} \right)$$

$$\omega_0^2 LC = L - R_L^2 C$$

$$\omega_0^2 LC + R_L^2 C = \frac{L}{C}$$

$$\therefore \frac{L}{C} = \omega_0^2 L^2 + R_L^2 \rightarrow \textcircled{2}$$

Substitute eq (2) in eq (1);

$$\therefore Y_0 = \frac{R_L}{(L/C)}$$

$$\therefore Y_0 = \frac{R_L C}{L}$$

$$\therefore Z_0 = \frac{L}{R_L C}$$

The impedance at resonance, i.e. " $\frac{L}{R_L C}$ " is called "dynamic impedance" of the ckt.

Definitions

1. Voltage magnification = $\frac{V_L}{V} = \frac{V_C}{V}$
(at Resonance)
2. Current magnification = $\frac{I_L}{I} = \frac{I_C}{I}$
(at Resonance)

- * Series Resonant ckt is also called as "SELECTOR CKT"
- * Parallel Resonant ckt is also called as "REJECTOR CKT"